##### Constraining Electron Parallel Energy in Electrostatic Fields through the Anomalous Doppler Effect Induced by External Electromagnetic Waves

A. Author , B. Author\*, C. Author

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**Abstract**

In this study, the interaction between free electrons and Electromagnetic Waves (EMW) in the presence of magnetic and electrostatic fields is simulated using the Volume-Preserving algorithm. We found that when the electric field of the EMW (which includes a left-hand polarization component) exceeds a critical value, it can continuously transfer the electron’s parallel energy into rotational energy via the Anomalous Doppler Effect (ADE). This mechanism converts the work by the electric field along the background magnetic field into perpendicular kinetic energy. As a result, the electron’s parallel kinetic energy saturates, while its perpendicular kinetic energy increases continuously. The model based on energy, momentum, and angular momentum conservation is given to provide a clear explanation of the contribution of left-hand polarization to the Anomalous Doppler Effect. This general model can be directly applied to interpret the anomalous Doppler effect in electron-wave interaction. To saturate the parallel energy of runaway electrons in magnetically confined plasmas, we propose an innovative solution for runaway electron mitigation in Tokamaks by launching extraordinary waves from the high-field side with an energy flux of several watts per square meter.

## 1 Introduction

During the current ramp-up phase in a tokamak, disruptions or magnetohydrodynamic (MHD) instabilities can induce quasi-static toroidal electric field, which accelerates electrons to energies on the order of several tens of MeV. This acceleration occurs when the force from the quasi-static electric field exceeds the drag from radiation and collisions. The accelerated electrons, known as runaway electrons, can cause significant damage to the interior walls of the device, thereby reducing its operational lifespan. It is worth considering the possibility of converting the acceleration of electrons from quasi-static electric fields into rotational energy within the magnetic field. This conversion not only suppresses the energy of runaway electrons, mitigating their detrimental effects on the device, but also enhances discharge performance by reducing the consumption of ohmic field energy . The transport of parallel energy in electrons to rational energy primarily occurs through three mechanisms: electron avalanche processes[25], collisionless pitch-angle scattering[22], and the Anomalous Doppler effect(ADE) [15]. The first two methods to suppress runaway processes, such as gas injection [35] and enhancing magnetic perturbations[37], usually have side effects and alter the discharge environment. In contrast, the latter method can be a clean process, making it particularly attractive for investigation.

When electrons move in static magnetic fields and interact with external electromagnetic waves (EMW) with frequency and wave vector , they experience a scattering phenomenon under the resonant condition, where and .The result of this scattering is a transfer of momentum from parallel motion to rotational motion, a phenomenon known as the Anomalous Doppler Effect. This effect was thoroughly described in the classic works of Ginzburg and Frank [15, 10, 24]. Recently ADE has raised increasing attentions in space radiation[8], runaway electron instability[21] and materials science[27], it is believed that ADE can explain the phenomena as whistler turbulence in flare loops[8], the Electron Cyclotron Emission (ECE) with step-like structure in Tokamaks [23, 6, 5]and the microwave bursts during Edge Localized Modes(ELM)[12] . Additionally, ADE holds potential for suppressing runaway electron energy in Tokamak discharges. This has been demonstrated by F. Santini, who found that high-energy runaways can be significantly reduced through ADE during lower hybrid wave heating in the Frascati Tokamak[26]. However, it is important to note that the high power of lower hybrid waves also causes an increase in nonthermal electrons through Landau resonance. This, in turn, leads to an increase in runaway electrons after shutdown of lower hybrid waves, which is a side effect of using lower hybrid waves to suppress runaway electrons.

Despite the Anomalous Doppler Effect (ADE) having been described in numerous experiments and theories[7, 4, 17, 28, 30, 13, 19, 14, 18], previous descriptions of the single-particle ADE have mainly relied on either quantum theory or classical theory, disregarding the accelerating electrostatic field[16, 26, 3]. Further exploration of the ADE in the presence of electrostatic fields is essential for understanding the physics of pitch-angle scattering of runaway electrons by electromagnetic waves in Tokamak discharges. Moreover, the nature of ADE in classical electrodynamics fields remains unclear due to the complexity of formulas in previous classical theories. For instance, it is still difficult to understand why parallel kinetic energy can convert to transverse internal energy during ADE resonance, or what kind of ElectroMagnetic Waves (EMW) can trigger ADE. These questions require further investigation.

This paper directly simulates the full orbit electron motion in uniform magnetic fields with accelerating electrostatic fields and electromagnetic fields using the Volume-Preserving Algorithm [38]. Compared to conventional algorithms,such as Boris, Volume-Preserving Algorithm can acquire long-term accuracy and conservativeness via a systematic splitting method, which is an excellent method for nonlinear electron dynamic simulation. To directly demonstrate the phenomenon of the ADE, we place an electron in a uniform magnetic field along with an electrostatic field that is oriented in the opposite direction to . As a result, the electron can be accelerated parallel to . Additionally, during the simulation, the slow electromagnetic wave with its phase velocity less than that of light in vacuum is set up as an induced wave. This induced wave allows us to directly observe the effects when the electron’s velocity reaches the resonant condition for the ADE. We explore the resonance with three types of polarization waves including linear polarization, left-hand circular polarization, and right-hand circular polarization in the simulation. We find only the wave with left-hand circular polarization can have ADE for runaway electrons. The critical energy of waves when the electron parallel velocity can be constrained and consistently transport the parallel energy from electrostatic field to transverse rotation energy is also found in the simulation. Furthermore, the self-consistency between the quantum theory and direct simulation about ADE is examined. Finally, after analyzing dispersion, polarization, and resonant moment, we have determined that the extraordinary wave is the most suitable for triggering ADE in plasma. Based on these findings, we propose a promising method for controlling runaway electrons.

## 2 Quantum theory of ADE

This extraordinary phenomenon has already been discussed from energy conservation in the works of V.L.Ginzburg[14], I. Tamm[32] , Nezlin[24],and I.M.Frank[11]. Here, we give an analysis based on the conservation of angular momentum.

As shown in Fig. 1, when charged particles move through a medium faster than the speed of light in the medium, induced currents are generated. These currents stimulate secondary waves that interfere with the electromagnetic field of the particles in motion, thus giving rise to Cherenkov radiation. The direction of Cherenkov radiation can only be along the Cherenkov radiation angle , where is the speed of light in the medium and v is the velocity of the charged particles. Now, let us substitute the charged particle with a system that has internal energy (for example an oscillator or a cyclotron electron in a magnetic field). The system moving faster than the speed of light emits photons with angular frequency and wavevector at direction , where the direction of the photon performs does not depend on the interference of the secondary waves and can be in any direction, as shown in Fig. 2. According to energy conservation and momentum conservation, we have :

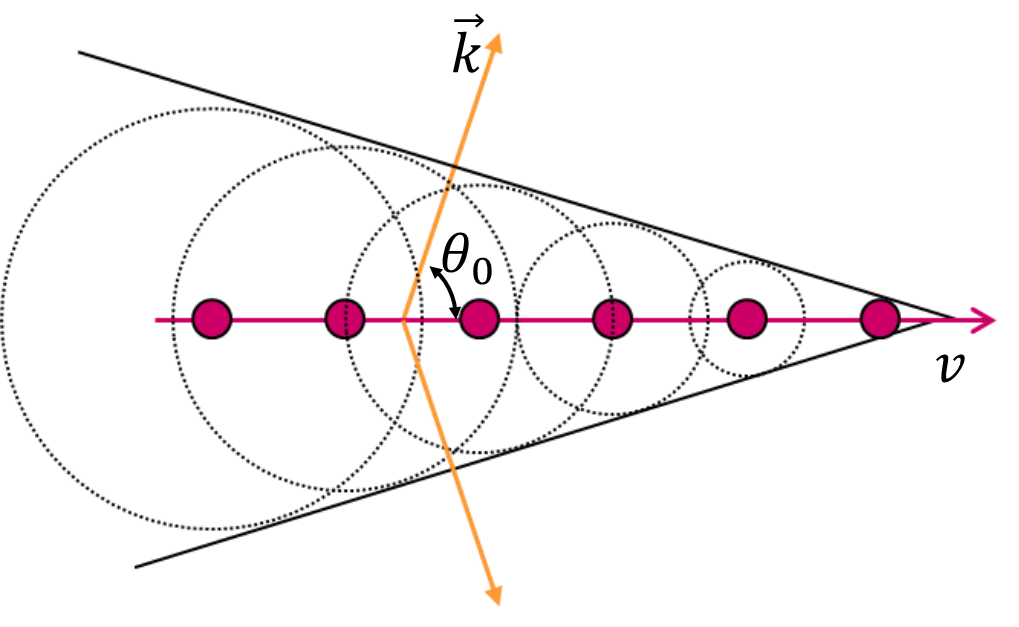


Figure 1: Schematic diagram of Cherenkov Radiation;the red point is the snapshot of the electron at different times

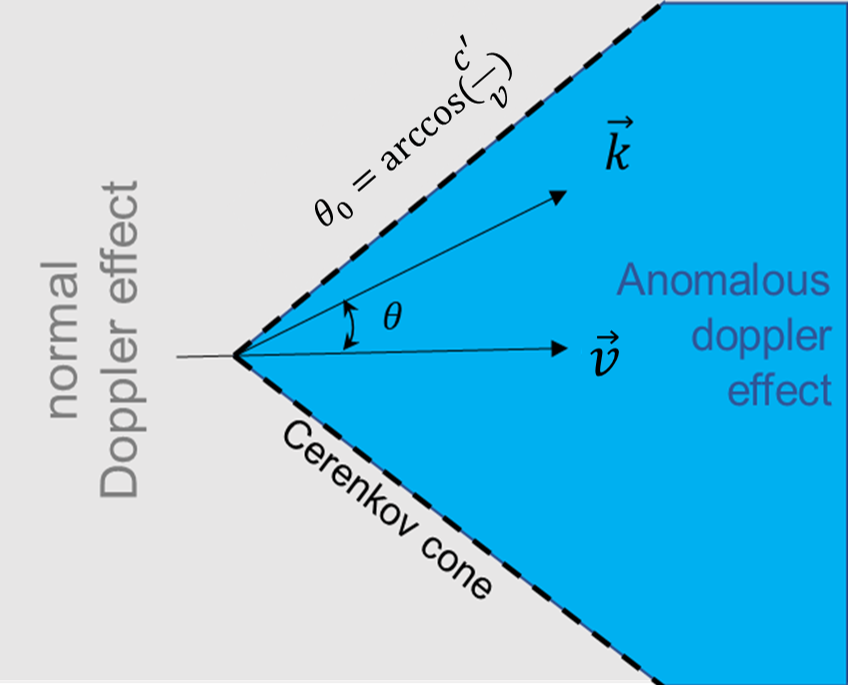


Figure 2: The region of ADE ,Normal Doppler Effect (NDE) and Cerenkov cone,where the blue region refers to ADE, the gray region refers to NDE, and while the dot line refers to the Cerenkov cone[27]

(1)

(2)

In the above, and represent the kinetic energy and internal energy of the system before emitting a photon, and and represent the energy of the system after emitting a photon. Considering that the photons energy is far less than the initial kinetic energy ,the losses of kinetic energy after emitting a photon can be expressed as , where is the velocity of the system before emitting a photon, and also . According to Eq. (1-2),we have :

(3)

Here, . While the system velocity is greater than the speed of light in the medium,according to the sign of , we can divide radiation into three regions, as shown in Fig. 2 : For , , The system produces photons by consuming its own internal and kinetic energy, this region refers to the Normal Doppler Effect (NDE); For , , this region is referred to the Anomalous Doppler Effect (ADE), where the system gains internal energy after emitting photons. It means the loss of kinetic energy is converted to photons and the system’s internal energy; For , , the loss of kinetic energy by the system is completely converted into photon energy; this line refers to the Cerenkov Effect(CE). All three effects are possible when the system velocity exceeds the speed of light . While the system velocity is less than the speed of light, only NDE exists. As we can see , we can judge the type of phenomena based on the change of internal energy after emitting photons.

Now consider the particular case in which the system is an electron moving freely with velocity in the direction of thee xternal magnetic field and having small transverse velocity component , then the kinetic energy of the cyclotron electron could be expressed as and the internal energy represented as , which is equal to the rotational energy of the electron as shown in Fig. 3. Here, the referred to is the Lorentz factor. In order to calculate the change of rotational energy, we should combine with another conservation: The angular momentum conservation, where:

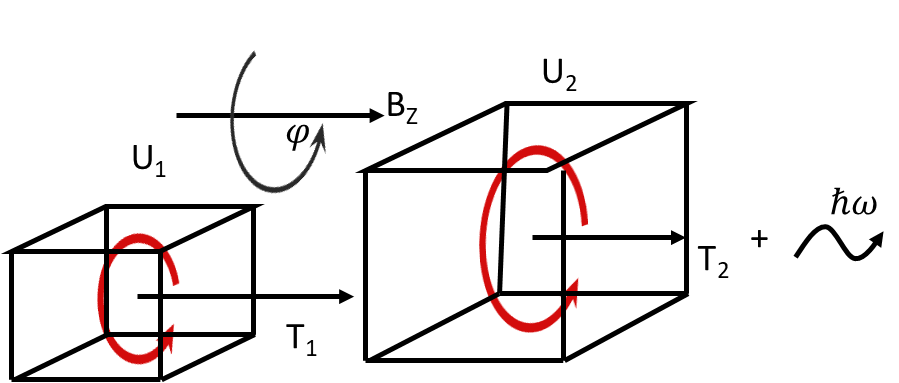


Figure 3: Schematic diagram of electron cyclotron system; here, , .

(4)

Here, we consider that the emitted photon contains angular momentum with and the angular momentum of cyclotron electron before and after emitting photon is and , since the magnetic field is along the z direction, should be equal to . From quantum theory, the electron wave in a static magnetic field can be expressed as:

(5)

In the above, the term is a normalized coefficient, A is the vector potential and s is the position. The intrinsic equation of angular momentum in the z-direction is :

(6)

The eigenvalue of is:

(7)

Consider ,and .We have :

(8)

where is the electron rest mass, is the Lorentz factor and is the electron cyclotron frequency in the rest frame. The variation in the angular momentum of the electron along z is :

(9)

In the above, m is the number of photon’s angular momentum in z direction. We the have :

(10)

According to Eq. (3) and Eq. (10), we have :

(11)

or

(12)

Here, represents the loss of kinetic energy , represents the energy of the photon, and represents the change in the electron cyclotron energy (also called the change in internal energy). There are three situations for :

1. When , , the cyclotron electron internal energy decreases after emitting a photon, and the emitted photon will have right-hand circular polarization with angular momentum to maintain angular momentum conservation. This process is called the Normal Doppler Effect .

2. For , , the Cherenkov Effect occurs, where the emitted photon does not cause any change in the internal energy of the cyclotron electron.

3. When , , the Anomalous Doppler Effect (ADE) occurs,resulting in an increase in the internal energy of the cyclotron electron and the emission of left-hand circular polarization with angular momentum .

As a result, the resonant condition is strongly associated with the wave’s angular momentum. For a plane wave, the wave angular moment number includes only . While for , it indicates that the resonant wave possesses a helicon structure. Based on the above discussion,there are three kinds of resonance for a system when electron moves along the uniform magnetic field with velocity under external EMW ;the resonant frequencies are Normal Doppler frequency(), Cerenkov frequency(), and anomalous Doppler frequency(). We only include assuming they are the most dominant resonances. These frequencies are respectively :

(13)

(14)

(15)

where is the angle between and . These equations are quite common resonant conditions for the kinetic equation of plasma ,what is intriguing is how the motion of electrons differs under various resonant conditions with an electrostatic field. However, this aspect has been far less researched in recent years.

## 3 Simulation framework

Considering uniform magnetic field in the z-direction, the electron is accelerated in the field by an electrostatic field that has

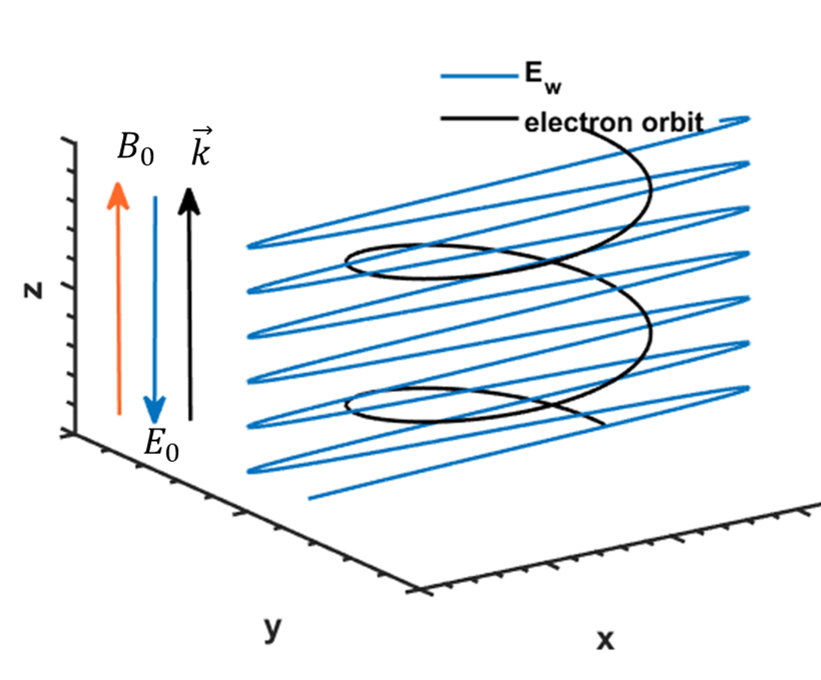


Figure 4: Schematic of simulation setup

opposite direction to as shown in Fig. 4. To analyze the dynamics of electrons during interactions with an electromagnetic field, we establish a plane EMW characterized by frequency and wavevector . The angle between and is denoted as ; the equation of motion of the electron is :

(16)

Here, and are the sum of the static field and EMW field. By using the VPA algorithm[38, 33, 34], the discrete structure of te differential Eq. (16) can be written as :

(17)

The operator Cay(A) denotes the Cayley transform of matrix A[38]; the dimensionless magnetic matrix is :

(18)

The dimensionless parameters are momentum , magnetic field , total electric field , time step , and position respectively, where the is the electron cyclotron period and .

We start with a linearly polarized wave with direction . To reduce the simulation time, is set to T . The wave angle frequency is , where . The wavevector is parallel to and . The amplitude of the electric field of the electromagnetic wave is and the electrostatic field is . All these parameters are set without reality considerations ,merely for rapid simulation . The time step always satisfies to ensure the accuracy of the simulation. The electron starts off stationary, but gradually gains speed .

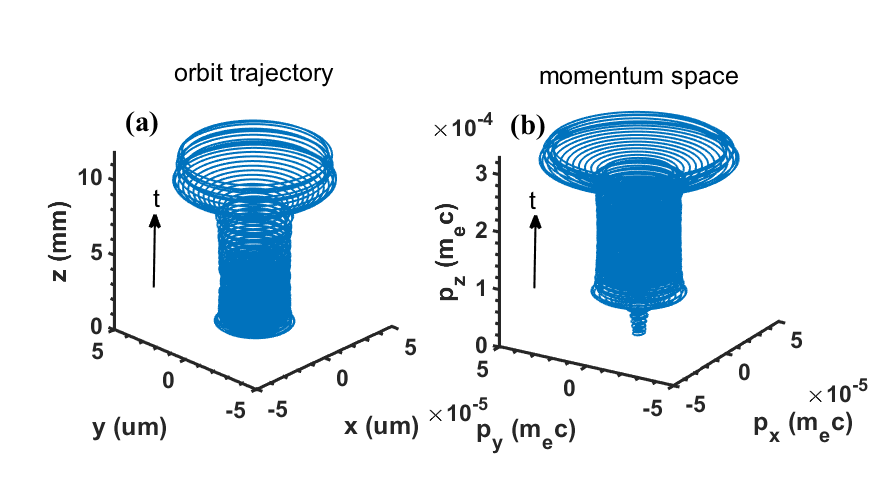


Figure 5: (a) Orbit trajectory of electron motion (b) Momentum phase space of electron motion

Simultaneously, the resonant frequency increases according to Eq. (13-15). Figure 5 shows the evolutions of the electron’s orbit and velocity phase during acceleration. The details of the electron’s motion are given in Fig. 6. The resonant frequencies increase simultaneously (Fig. 6(a)) during electron acceleration in the electrostatic field (Fig. 6 (b)). When the Normal Doppler frequency is equals to that of the induced wave at about , the perpendicular velocity (or rotational velocity) rapidly increases (Fig. 6(d)). The parallel velocity caused by EMW also increases simultaneously , as shown in Fig. 6(c), which can be calculated from where is the parallel velocity in EMW and electrostatic field while is only the result of electrostatic field . This phenomenon represents the NDE where the resonant velocity is "subluminal". The absorption of induced waves by the cyclotron electron causes an increase of velocity in the parallel and perpendicular directions which can be viewed as a reverse process for a system emitting photon as discussed previously (The NDE process is widely employed for current drive [2] and plasma heating [20] in tokamaks , although it is generally believed that the current drive by electromagnetic waves follows the Fisch mechanism [9] due to the limited toroidal momentum injected by the waves). However, the resonant condition is quickly broken as the parallel velocity keeps increasing until it reaches , at which point the ADE starts to come into view. When time reaches , the system starts to resonate with the induced wave through ADE, where the as shown in Fig. 6(a), the parallel velocity begins to scatter to the perpendicular direction as we can see from the decrease of and the increase of shown in Fig. 5(c) and Fig. 6(d). The resonant condition quickly disappears as the parallel velocity surpasses the resonant region. To determine which type of electromagnetic wave is responsible for the NDE and ADE separately, we separate the linear polarized wave into left-handed circular polarization and right-handed circular polarization. We observe that the right-hand circular polarized wave is responsible for the NDE, while the left-hand circular polarized wave induces the ADE, as illustrated in Fig. 6(e) which agree well with our previous analysis. We can understand this phenomenon through the conservation of angular momentum and linear momentum: Because electrons have a right-hand spin orbital angular momentum in a magnetic field, when the electron absorbs right-handed electromagnetic waves propagating in the parallel direction, according to the conservation of momentum and angular momentum, the parallel momentum and rotational energy of the electron will also increase; this process corresponds to NDE. When the electron emits left-hand circular electromagnetic waves propagating in the parallel direction, the conservation of momentum results in a decrease in the electron’s parallel momentum, while the conservation of angular momentum requires an increase in the electron’s rotational energy;this process corresponds to ADE.It is worth noting that there is no response during the Cerenkov resonance for the electromagnetic wave (EMW), as the Cerenkov effect primarily pertains to electrostatic waves.The ADE mechanism from classical theory is analyzed in the Appendix.

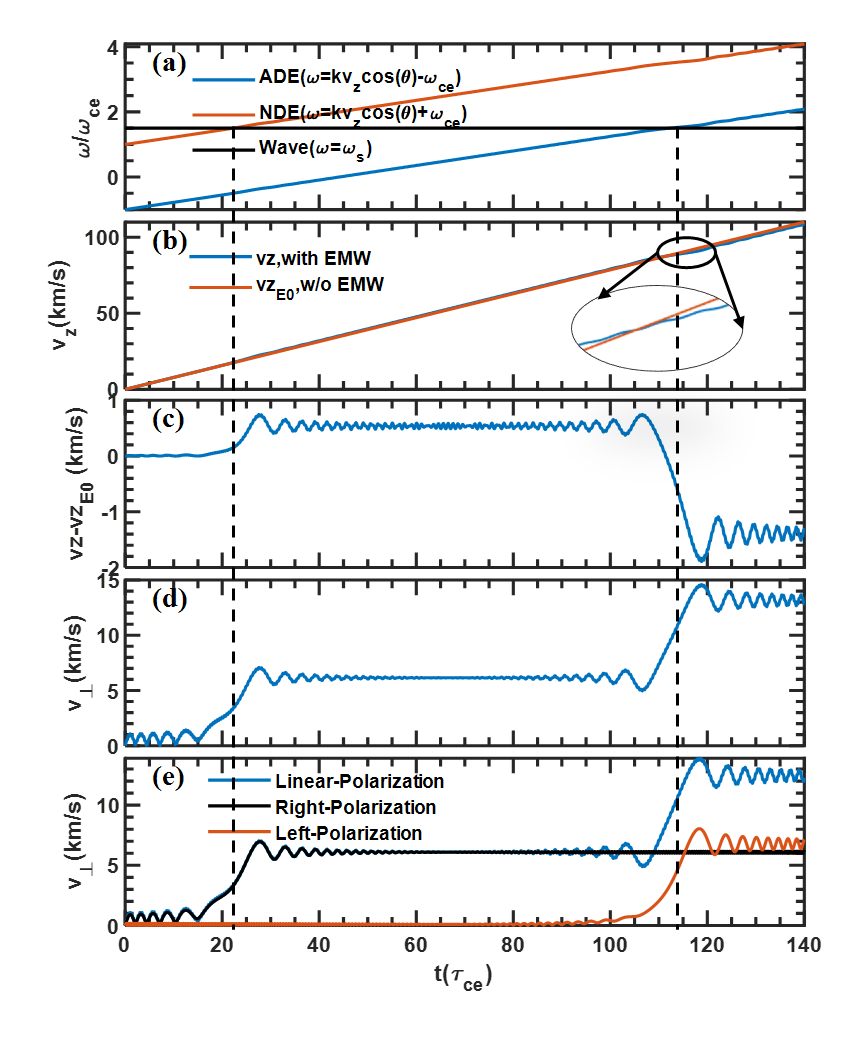


Figure 6: Kinetic evolution of electrons in a magnetic field with electromagnetic wave during acceleration. (a)Wave frequencies of anomalous Doppler frequency, normal Doppler frequency, and source wave frequency,(b)Translational velocity with and without electromagnetic wave ,(c)The change of parallel velocity caused by electromagnetic wave, (d)Cyclotron velocity , and (e)Interaction with Linear ,Right-hand circular ,Left-hand circular polarization .

## 4 Critical Trapping Threshold of ADE

We found that the ADE is equivalent to an effective damping force, hindering the electron acceleration process. Therefore, if the intensity of the electromagnetic wave is increased, it is theoretically possible to achieve a balance with the electrostatic field force, causing the electron to no longer be accelerated by the electrostatic field. Next,we will demonstrate the existence of this equilibrium by varying the electromagnetic wave field intensity .

Let us consider an electromagnetic wave with only a left-hand polarized circular component, where the wave-vector k and frequency of the EMW are set as and . Here, and is parallel to the static magnetic field. The electrostatic field and the static magnetic field are set as and respectively. As shown in Fig. 7, when we increase the energy of the electromagnetic wave, the parallel velocity will be trapped in the resonant condition and ceases to increase while the perpendicular velocity increases continuously when the ratio exceeds a certain threshold. The trapped electron’s orbit and momentum are shown in Fig. 8.

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Figure 7: Time trace of velocity under different ratios of . (a)vertical velocity (b)parallel velocity (c) Zoom in parallel velocity.

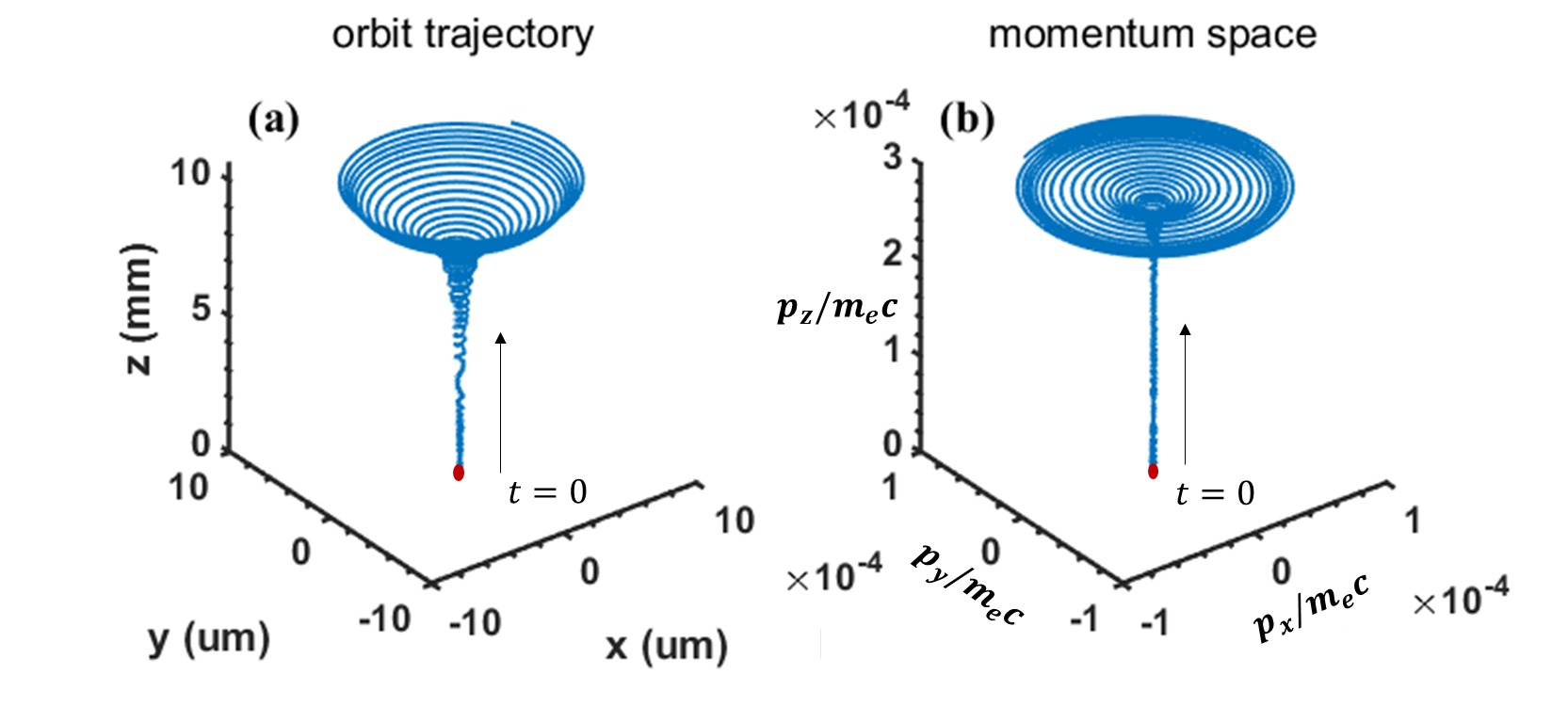


Figure 8: The electron’s orbit and momentum with trapped parallel energy

The threshold field E can be obtained by modifying the electromagnetic wave intensity using a dichotomy control approach based on the final parallel velocity with a sufficiently long time. The critical ratios of with dimensionless parameters are shown in Fig. 9 (here, the angles between and wavevector are ). As observed, a larger angle between k and results in a less critical ratio for trapping parallel energy. This possibility could be attributed to the increase in the electric field along the parallel direction as the angle increases, which leads to an increase in the retarding force along that direction. Furthermore, keeping the magnetic field strength and the wave frequency constant, a decrease in the wavevector results in an increase in the critical field strength. This behavior arises because a lower corresponds to a higher resonant velocity. Consequently, the power imparted by the static electric field, given by , also increases. As a result, a higher electromagnetic wave intensity is required to facilitate timely conversion of the parallel energy.

It is apparent that with only small intensity of left-hand circular polarized wave we can halt the increase of the electron’s parallel momentum and transfer energy from the electrostatic field to rotational energy by . For example, in tokamaks the toroidal electric field is about ; consequently, the threshold electric field for a left-hand circular polarized wave to trap parallel energy is about in the plasma (the corresponding energy flux is about for and with .

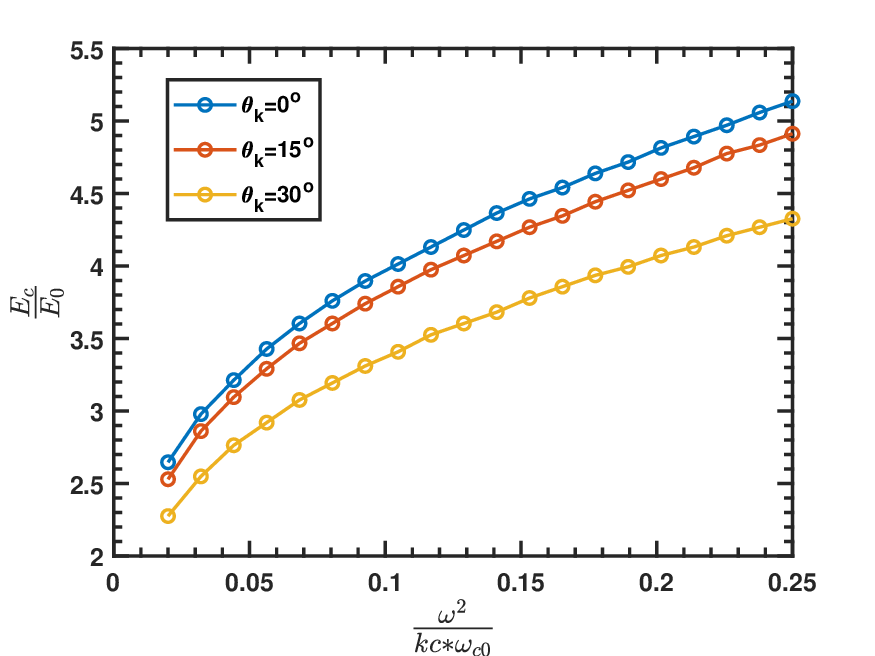


Figure 9: ,where , , and the refractive index ranges from 4 to 50.

## 5 Simulation with Tokamak parameters

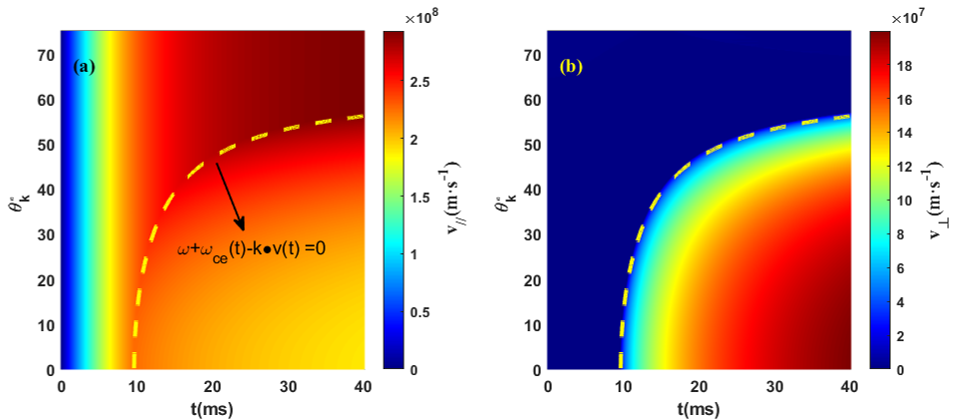
In order to check the validity in high magnetic fields and the angle dependence of ADE, we choose a uniform field of and , which are typical parameters of a Tokamak during startup[21]. For a plane left-hand circularly polarized wave, , and ,the energy flux of the wave is about . Using the powerful parallelism of a supercomputer, we can simultaneously calculate the interaction of 500 electrons with a left-handed circularly polarized wave at 500 different incident angles in the range of 0 to 90 degrees, demonstrating the dependence of the ADE effect on the angle. The time step in the simulation is set at to ensure the convergence of the results. The results are shown in Fig.10, where the dashed yellow line represents the earliest resonant time that satisfies the relation . It is evident that as the angle increases, the onset of resonance is delayed, and the requisite speed for resonance is heightened. Following the initiation of resonance, the rotational velocity witnesses a swift increase, whilst the parallel velocity gets ensnared within the resonant region, thereby ceasing any further increment. 

Figure 10.Time evolution of and by electromgnetic wave with different wave incident angle

These results indicate that even with a small energy of left-hand polarized waves, we can suppress runaway electrons and transfer the acceleration from the electrostatic field to rotational energy in Tokamak parameters.

## 6 Electromagnetic wave drives ADE in magnetized plasma

Runaway electrons present a significant threat to tokamak devices due to their potential to damage plasma-facing components (PFCs)[36] through high-energy impacts. Harnessing the ADE effect in tokamaks offers a promising approach to controlling runaway energy, provided three critical conditions are satisfied:

1. The existence of a left-hand polarized wave component when electrons move in the direction of the magnetic field;

2. The phase velocity of the electromagnetic wave must remain below the speed of light in vacuum.

3. The wave must possess sufficient energy to counteract the accelerating velocity of the electrons. Meeting these conditions enables the establishment of controllable resonant interactions essential for mitigating runaway electrons.

Consider waves in cold magnetized plasmas where ,from Appleton-Hartree formula(It is commonly believed that the cold plasma dispersion is generally adequate for calculating wave propagation when the phase velocity significantly exceeds the thermal velocity of electrons in the plasma.):

(19)

In the above, is the wave frequency and k represents the wavevector. Both and depend on the electron cyclotron frequency(), the plasma frequency () and the ion cyclotron frequency(). Here, is the angle between and the static magnetic field . The dispersion relationship can be illustrated in Fig. 6(a), where the blue color represents the wavevector , while the red color represents . The polarization component in direction of the wave can be expressed as[1]

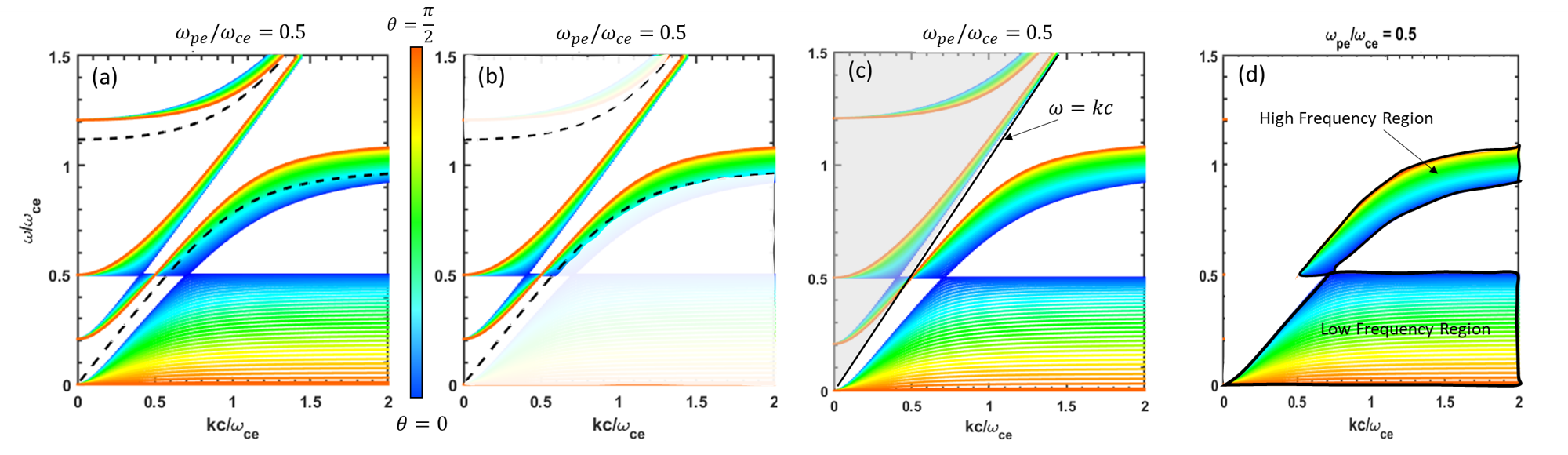
 (20)

Figure 11.Dispersion relationship in cold magnetized plasma. (a) Complete dispersion relationship in cold magnetized plasma. (b) Region dominated by left-hand polarization. (c) Region where the phase velocity is smaller than the speed of light in vacuum. (d) High-frequency and low-frequency regions.

The electric field of the wave can be written as . The left-hand polarized wave along z can be writen as ,while the right hand polarized wave is . The component of the left-hand polarized wave in E can be calculated by and for right-hand polarized wave. When B is positive, we have ; the main component of the wave is a right-hand polarized wave. Otherwise, it is predominantly a left-hand polarized wave. The black dashed line in Fig. 11(a) represents the boundary between two different types of waves where and the dispersion relationship in cold magnetized plasma.Therefor as shown in Fig. 11(b), between the pair of dashed lines, the wave predominantly exhibits left-hand polarization. Conversely, beyond these lines of demarcation, the wave is mainly right hand polarized. In Fig. 11(c), the black line refers to the waves in the vacuum. only waves below the black line have the ability to cause ADE, where the phase velocity . Finally, we found the wave for ADE as depicted in Fig. 11(d), low frequency region and high frequency region. The low frequency region comprises whistlers and magnetized electron plasma waves, where all electromagnetic waves mainly exhibit right-hand polarization. Conversely, in the high frequency region , extraordinary waves display left-hand polarization when is in close proximity to as shown in Fig. 11(b).

Consider the resonant condition

(21)

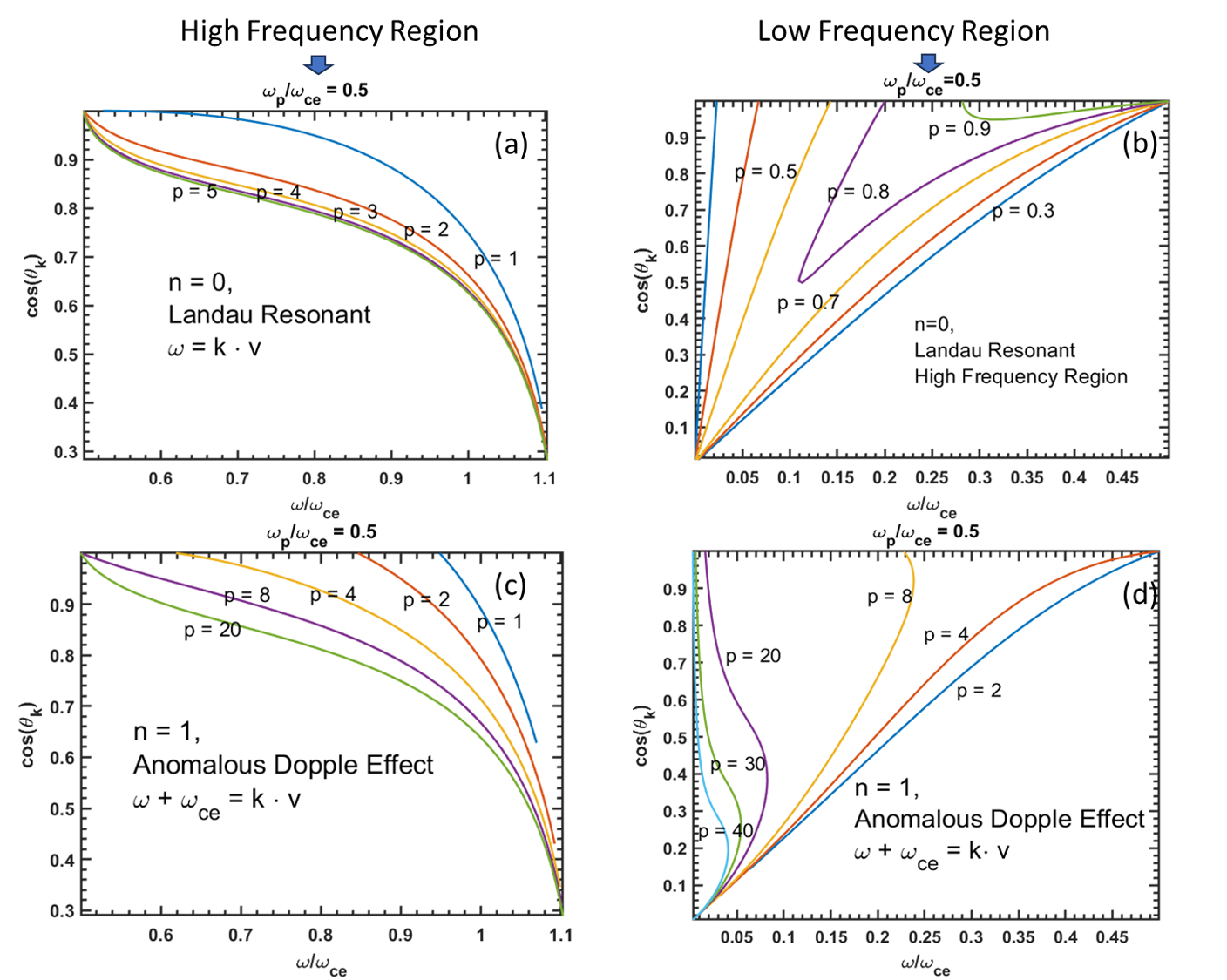
When the runaway electron’s momentum matches the resonant condition, it is possible to stimulate the intrinsic waves in plasma. To simplify the calculation here we assume the velocity along the static magnetic field and we choose for , and for Landau resonance. The term is disregarded in this context since NDE is not important in the high frequency region due to fact that it requires a right-hand polarized wave . By combining Eq. (19)and Eq. (21), we establish the relationship between the wave and resonance moment, as illustrated in Fig. 12. In the high frequency region , for , resonance curves with dimensionless resonance momentum greater than unity converge to the right bottom region shown in Fig. 12(c), which corresponds to the Extraordinary wave with frequency range . This explains why Extraordinary wave near the frequencies of are excited when runaway electron scattering occurs in magnetized devices [28, 12]. Furthermore, the dimensionless Landau resonant momentum for most of the low frequency region is greater than 1, as shown in Fig. 12(a). This suggests less attenuation of waves by background thermal electrons and makes wave formation in the low frequency region easier. For the high frequency region, when the energy of high-energy runaway electrons exceeds (reduced momentum ), the resonance curves of electromagnetic waves excited by the effect almost always pass through the left top region depicted in Fig. 12(d). This region corresponds closely to the whistler wave zone where whistler waves propagate parallel to the magnetic field. Therefore, in Tokamak experiments, the observation of whistler waves is typically associated with the detection of high energy electrons with energy exceeding [31]. The dimensionless Landau resonance momentum for low frequency region is less than unity, as depicted in Fig. 6(b). This 

Figure 12.In the high frequency region, the lower left boundary corresponds to the upper hybrid wave at different angles. In the low frequency region, the lower right boundary corresponds to the lower-hybrid wave at different angles. (a) (b)dimensionless moment of Landau resonance in the high frequency region and B. (c) (d)dimensionless moment of ADE in the high frequency and low frequency region,unit using for dimensionless :

indicates a higher degree of wave attenuation by background thermal electrons compared to high frequency region, making wave formation in region more challenging.

Based on the above discussion, electromagnetic waves in the high frequency region are more prone to exciting Anomalous Doppler Resonance due to polarization and damping, while waves in the low frequency region are better suited for heating background electrons through Landau resonance, such as lower-hybrid wave heating. Experiments show that runaway electrons can stimulate extraordinary waves with frequencies in the range to through and transport their parallel energy to rotational energy[12, 29]. It is natural to consider utilizing the reverse process by injecting extraordinary waves to suppress runaway electrons.

## 7 injection of extraordinary wave to suppress runaway electrons

The characteristic frequency with in a tokamak plasma is shown in Fig. 13 under conditions where the intensity on the axis toroidal is and the density along the minor radius has a profile given as , and r is the normalized minor radius.

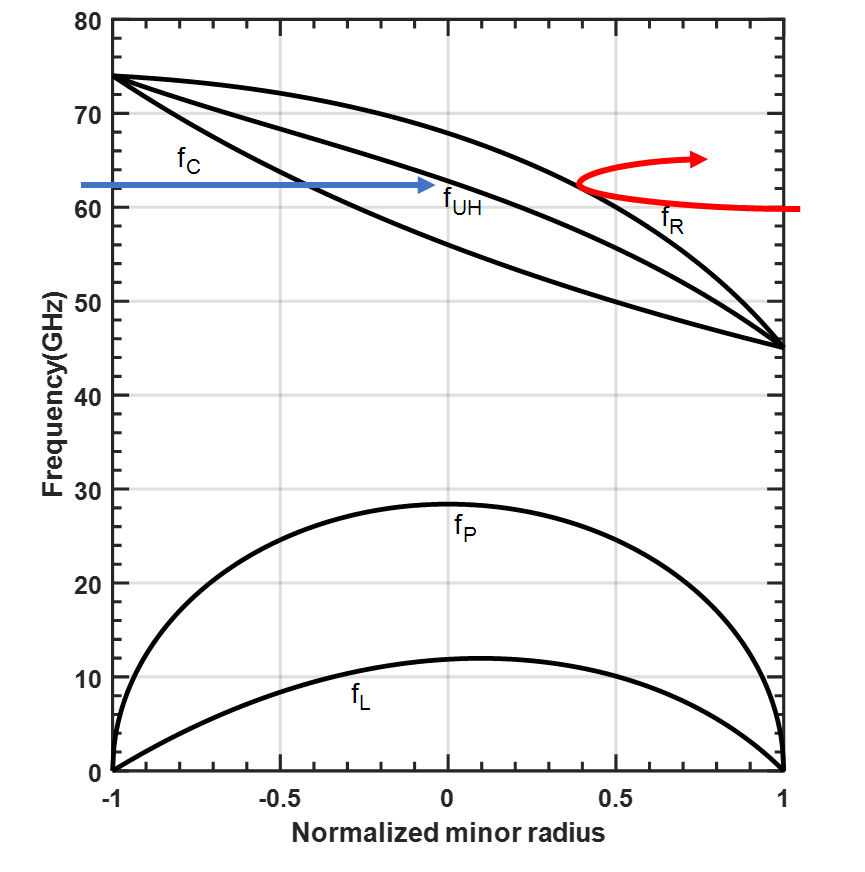


Figure 13: Characteristic frequency distribution in plasma. is the electron cyclotron frequency,is th upper hybrid frequency, is the plasma frequency and is the left-hand cutoff frequency, and is the right-hand cutoff frequency.

The X-Wave will slow down near the upper-hybrid frequency layer and reflect at the right hand cut-off frequency layer. Therefore, it is necessary to inject the extraordinary wave from the high field side of the tokamak, whereas it will be reflected at the right cut-off frequency layer when injected from the lower-field side. Different frequencies correspond to different positions. We can adjust the frequency of the extraordinary wave to align with the region where runaway events are more likely to occur (such as in the core of the tokamak). Since the power requirement for trapping runaway electrons is not higher than kilowatts, it is worth noting that precise frequency adjustments to match the core of the Tokamak based on real-time plasma density diagnosis is possible to achieve.

## 8 Summary

This study presents a simulation of the interaction between electrons and plane electromagnetic waves within uniform magnetic and electrostatic fields. The simulation demonstrates that the parallel velocity of electrons can be rapidly converted into rotational velocity during the ADE, with only left-hand circular polarization waves contributing to the ADE. When the electric field surpasses the critical field, the electromagnetic wave can capture the parallel momentum of electrons and continuously transfer energy from the parallel electrostatic field to rotational energy and resonance waves. The proportion between the loss of parallel energy and the augmentation in rotational energy aligns with both quantum theory and the simulation results. Furthermore, this study also delves into the fundamental physics of ADE from the basic energy-momentum-angular momentum conservation. We ascertain that the extraordinary wave, situated near the upper-hybrid resonance, emerges as the most apt candidate wave for mitigating runaway electrons. By injecting the extraordinary wave from the high-field side of the Tokamak, we can potentially inhibit the runaway electrons at the Tokamak core. It is possible to suppress runaway electrons at the core of the Tokamak based on real-time plasma density diagnosis by dynamically adjusting wave frequencies, where the field intensity needed to achieve this suppression is only a few orders of magnitude greater than the electrostatic field .

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## 9 Appendixes

Further investigation is necessary to comprehend why left-handed circular polarization is responsible for the ADE and what causes the scattering of the electron’s parallel momentum to rotational moment. In this Appendixes, our objective is to understand this phenomenon using classical electrodynamic theory.

### 9.1 EMW polarization

During the resonance with the induced wave, the rotation of the electromagnetic field and electron motion become synchronized. The wave frequency is equal to the electron cyclotron frequency in the frame of a gyro-center moving electron. The frequency and wavevector in the two frames, namely the laboratory frame and the moving gyro-center electron frame, can be directly linked through the Lorentz transformation [27]:

(22)

where and are wavevector and the frequency in the laboratory frame (the moving gyro-center electron frame) respectively. is the parallel velocity of the electron with normalized form being , is the Lorentz factor. where is unit tensor. From relationship (13),

(23)

(24)

Substitute with and combine with (14)(15), we get the frequency in laboratory frame

(25)

where representing the right-hand rotation for an electron in magnetic field. In the case of the NDE,it is essential to note that the resonant velocity, denoted as , must satisfy the condition . This condition ensures that both angular frequencies, and (where ), share the same sign. Consequently, it proves only waves exhibiting right-hand polarization can induce the NDE. In the case of the ADE, a different scenario arises. The resonant velocity is expressed as . For ADE to occur, the condition must be satisfied. Consequently, it is evident that and have opposite signs in Eq.(25), indicating that only left hand circular polarization waves are responsible for inducing the ADE .This feature has been rarely considered in previous theories of ADE. It is worth noting that the requirement of left-hand circular polarization for performs does not mean that the wave should be completely left-hand circularly polarized. As long as there exists a component of left-hand circular polarization in the direction of the cyclotron electron’s motion,it is possible to trigger the ADE , as observed with whistle waves in Tokamak systems [30].

### 9.2 Scattering

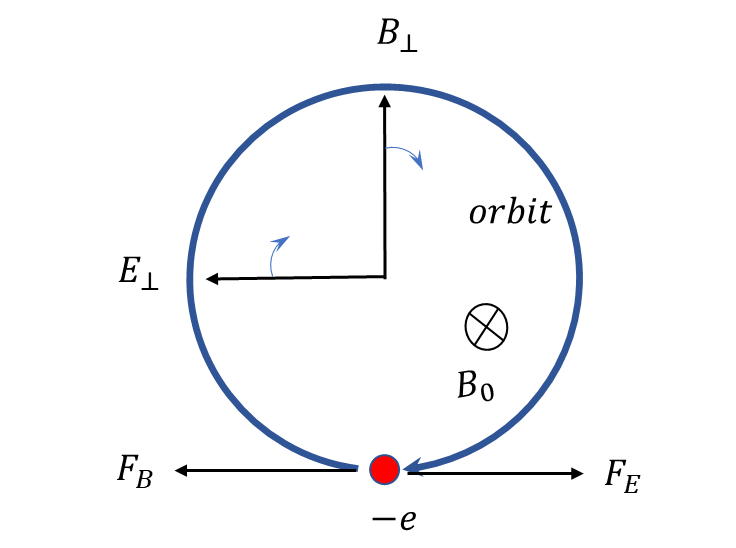


Figure 14: Scheme of electron movement in uniform magnetic field with EMW,here is static magnetic field , and are refer to electric field and magnetic field of EMW respectively. represents magnetic force of EMW, represents electric force of EMW

The scattering results shown in Fig.6 can be interpreted from classical electrodynamics as follows: Considering the circular polarization of EMW has the same frequency as during resonance in the moving gyro-center electron frame, as shown in Fig. 14, where the wavevector k parallel to the z-axis ,the perpendicular forces exerted by the magnetic field ((B)) and electric field ((E)) of EMW can be described as follows:

(26)

(27)

the forces and have opposite directions. The net force on electron in perpendicular direction is

(28)

where ,.The electron velocity has tendency to parallel to by changing the rotation phase of electron during resonant with EMW, as shown in Fig. 10.The force from EMW paralleled to z is:

(29)

the total work by EMW is

(30)

For NDE, , the rotational velocity tend to parallel to under external EMW,then ,,EMW will heat and exert a parallel force, causing the electron to accelerate along the z-axis.

For the Cerenkov Effect(CE), where and , the absence of a vertical force from EMW results in a random phase angle between and ,so the average of , , which means that the EMW doesn’t take part in Cerenkov effect. The longitudinal wave is more relevant to Cerenkov resonant, which also known as landau resonant, where the electron rides on the crest of the wave, being exposed to a force thus continually absorbs energy from the wave until its velocity increases and drops out of phase with the wave. such as lower hybrid waves in plasma[26] .

For ADE, , the rotational velocity tend to parallel to under external EMW,then ,.As a result,the electron will convert its kinetic energy into rotational energy and emit EMW. It is interesting to note that the magnetic field of the EMW performs positive work on the electron in the perpendicular direction and negative work in the parallel direction. The behavior of magnetic field in the EMW acts like a bridge connecting the parallel and perpendicular directions, transferring energy from kinetic energy to rotational internal energy. This fundamental physics process helps us understand why parallel kinetic energy can be converted into rotational energy.

### 9.3 Consistency between quantum theory and simulation

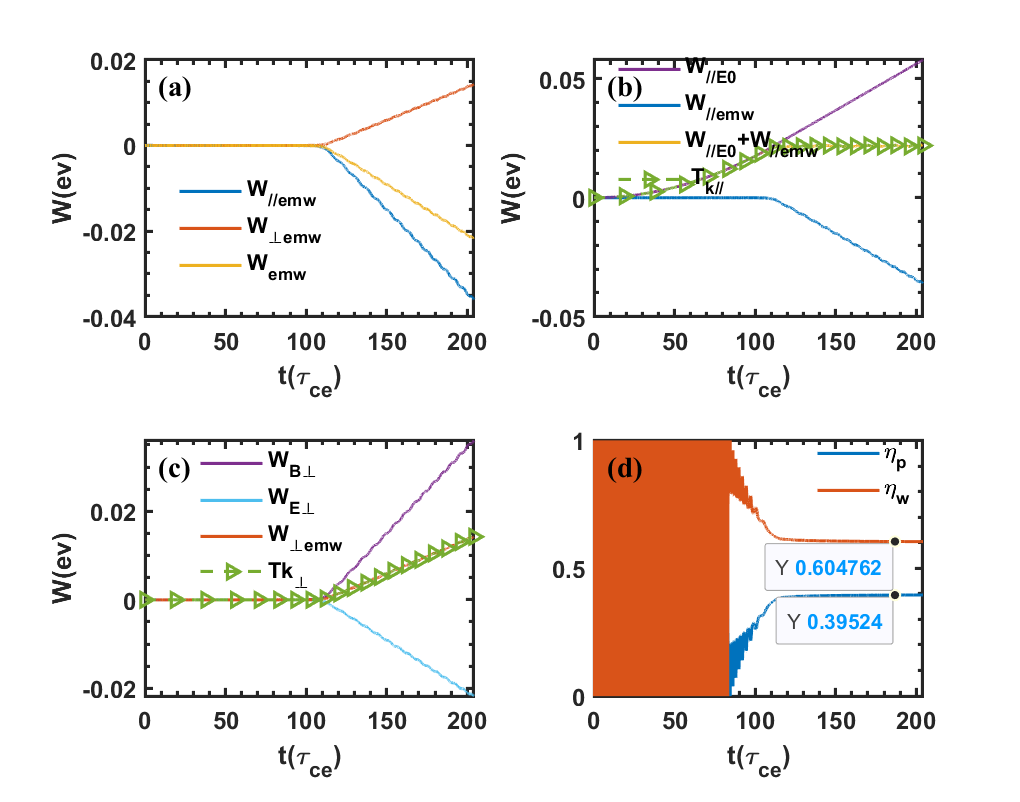
This section focuses on analyzing the energy flow during the Anomalous Doppler Resonance of electron excited by a left-hand circularly polarized wave with wave vector k parallel to the magnetic field . In the vertical direction, the electron is only subjected to work by the electromagnetic wave, which can be divided into two components: , representing the work done by the magnetic field of the electromagnetic wave, and , representing the work done by the electric field of the electromagnetic wave. Here, and . The total work of the EMW in the vertical direction is . In the parallel direction, both the electrostatic field and the electromagnetic wave can exert work on the electron. This includes , representing the work done by the electrostatic field, and , representing the work done by the electromagnetic wave. Specifically, and ,where is Eq.(29). The total work in the parallel direction is given by . All this work can be calculated numerically. Here, we use the same setup as above with , where the electron can be trapped by the EMW. Numerical results are shown in Fig. 15, where and are rotational energy and parallel kinetic energy of electron respectively. Fig. 15(a) shows that EMW performs negative work on the electron in the parallel direction while positive work in the vertical direction,the total work of EMW on electron is negative ,which means the electron transfers energy to EMW . In the parallel direction, as shown in Fig. 15(b) ,the electromagnetic wave performs negative work, which is balanced by the positive work from the electrostatic field. This equilibrium causes the electron’s velocity to stop increasing.In the vertical direction(Fig. 15(c)),the magnetic field of the EMW performs positive work, while the electric field of the EMW performs negative work. The total work in the vertical direction done by the EMW is equal to the change in kinetic energy of the electron, denoted as . This observation aligns well with the previous analysis(9.2). 

Figure 15.Energy analysis of EMW and electrostatic on the resonant electron during ADR(Anomalous Doppler Resonance). (a) EMW’s energy analysis for both direction (b) parallel direction. (c)vertical direction (d) conversion ratio of electrostatic work to EMW () and to internal energy of electron ()

Let’s consider the energy transformation coefficient.The decreases energy of the electron in the parallel direction is , while the increase energy of the electron in the vertical direction is ,the increased energy of EMW is .The ratio of the loss of parallel energy of the electron to the vertical energy is represented by , while to EMW is .The numerical results of and are depicted in Fig. 15(d) and .According to quantum theory, the decreases energy of the electron in the parallel direction is . The increase in rotational energy is , where we choose for the anomalous Doppler effect at here, the photon energy is . According to the simulation setup with the following parameters: , , , and , the resonant velocity is approximately . Consequently, the ratio of , and the ratio of . These values align well with classical electrodynamics calculations as shown in Fig. 15(d).

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